### TPMN 2019/2020 : Solving the Schr"odinger

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<sup>e</sup> Aim of this course : Solving the Sch<sup>~</sup>odinger equation by using a computer

 $H\Psi(r) = s\Psi(r)$ 

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 $H\Psi(r) = s\Psi(r)$ 

#### <sup>e</sup> Three kind of terms

 $^{\bullet}$  H, the Hamiltonian operator  $\rightarrow$  it describes the quantum system ; in general it is known

 $\Psi(r)$ , the wavefunction of the system ; we want to compute it

es, the total energy of the system ; we want to compute it

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<sup>e</sup> **Problem :** the Schr odinger equation is an eigenvalue problem  $\rightarrow$  to get  $\Psi(\mathbf{r})$  we need s

we to get s we need  $\Psi(r)!$ 

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s

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- <sup>e</sup> From a numerical point of view, there are different ways to solve this problem
  - <sup>e</sup> Tomorrow, we will see a quite general method : the Finite Difference Method
  - Next week, R. Hertel will present another possible way to proceed : the Finite
     Element Method

<sup>e</sup> The problem to solve : Free particule in a box (1D)



<sup>e</sup> The problem to solve : Free particule in a box (1D)



$$-\frac{1}{2}\frac{d^2\Psi}{dx} = s\Psi$$

### The Numerov algorithm <sup>e</sup> The problem to solve : Free particule in a box (1D)



### The Numerov algorithm <sup>e</sup> The problem to solve : Free particule in a box (1D)



This differential equation belongs to the general kind of 2<sup>nd</sup> order linear differential equation

$$\frac{d_2\Psi}{dx} + Q(x)\Psi(x) = S(x)$$

where Q(x) and S(x) are continuous functions on a domain [*a*, *b*]. The equation is to be solved as a **boundary value problem**, *i.e.*,  $\Psi(a)$  and  $\Psi(b)$  are given.

$$\frac{d_2\Psi}{dx} + Q(x)\Psi(x) = S(x)$$

Depending of the functions Q(x) and S(x), the Numerov algorithm can be used to solve

Eigenvalue problem: Q(x) f = 0 and s(x) = 0 The Schrödinger equation Ex.: Hydrogen atom  $-\frac{k^2}{2m} = \nabla_{2\psi}(r - \frac{1}{r} \frac{Ze^2}{r} \psi(r) = s\psi(r) \xrightarrow{symential}}{\psi_{nlm}(r)} = \gamma_{lm}^{u_{nl}(r)}(\theta, \phi)$   $\gamma_{lm}(\theta, \phi)$  are the spherical harmonics and the function u(r) is given by  $2^{nd}$  order differential equation  $\frac{d^2u}{dr} = -Q(r)u(r)$  with  $Q(r) = 2s + \frac{2Z}{r} = \frac{L(l)}{r}$ 

$$\frac{d_2\Psi}{dx} + Q(x)\Psi(x) = S(x)$$

Depending of the functions Q(x) and S(x), the Numerov algorithm can be used to solve

Linear system problem: Q(x) = 0The 1D Poisson equation Ex. : Hartree potentiel in spherical symmetry  $V_{Hartre}$  (r) =  $\frac{1}{\frac{e^2}{4\pi s^0}}$  d<sup>3</sup><sub>r</sub>  $j\frac{\rho(.)}{|.-.|}d^2U$  symmetry  $= -4\pi r \rho(r)$ Hartree dr <sup>2</sup>  $U_{Hartree}(r) = rV_{Hartree}(r)$ d<sub>2</sub> U dµartree  $S(r) = 4\pi r \rho(r)$ ・ロト ・ 日本 ・ モト ・ モト ・ モー O(x) = 0

### The Numerov algorithm <sup>e</sup> The problem to solve : Free particule in a box (1D)





#### The Numerov algorithm The problem to solve : Free particule in a box (1D) $\frac{d_2\Psi}{dx} + Q(x)\Psi(x) = S(x)$



$$\frac{1}{2}\frac{d^{2}\Psi}{dx^{2}} = s\Psi$$
$$\frac{d_{2}\Psi}{dx^{2}} + 2s\Psi(x) = 0$$

$$\frac{1}{2} \frac{1}{2} \frac{1}$$

#### The Numerov algorithm The problem to solve : Free particule in a box (1D) $\frac{d_2\Psi}{dx} + Q(x)\Psi(x) = S(x)$



$$\frac{1}{2}\frac{d^{2}\Psi}{dx^{2}} = s\Psi$$

$$\frac{d_2\Psi}{dx^2} + 2s\Psi(x) = 0$$
$$Q(x) = 2s$$
$$S(x) = 0$$

$$-\frac{1}{d^2 \psi x^2} + \mathbf{V}(\mathbf{x}) \Psi(\mathbf{x}) = \mathbf{s} \Psi$$

#### The Numerov algorithm The problem to solve : Free particule in a box (1D) $\frac{d_2\Psi}{dx} + Q(x)\Psi(x) = S(x)$



$$\frac{1}{2}\frac{d^{2}\Psi}{dx^{2}} = S\Psi$$

$$\frac{d_2\Psi}{dx^2} + 2s\Psi(x) = 0$$
$$Q(x) = 2s$$
$$S(x) = 0$$

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$$-\frac{1}{d^{2}\Psi x^{2}} + V(x)\Psi(x) =$$
  
s\Phi = 0  

$$\frac{d^{2}\Psi}{dx} + 2(s - V(x))\Psi(x) = 0$$
  

$$Q(x) = 2(s - V(x))$$
  

$$S(x) = 0$$

<sup>e</sup> We consider a grid, step  $\Delta$ ,



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<sup>e</sup> We consider a grid, step  $\Delta$ , <sup>w</sup> We resort to Taylor series to express  $\Psi(x + \Delta)$  and  $\Psi(x - \Delta)$   $\Psi(x + \Delta) = \Psi(x) + \Delta \frac{d\Psi(x)}{dx} + \frac{\Delta^2 d\Psi(x)}{dx} + \frac{\Phi^2 d\Psi(x)}{dx^2} + \frac{\Phi^2 d\Psi(x)}{dx^2} + O(\Delta)^5$  $\Psi(x - \Delta) = \Psi(x) - \Delta \frac{d\Psi(x)}{dx^5} + \frac{\Delta^2 d\Psi(x)}{dx^4} - \frac{\Delta^3}{6} + \frac{\Delta^4 \Psi(x)}{dx^4} - O(\Delta)^5$ 

 $\Psi(x)$ 

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dx <sup>2</sup> dx <sup>6</sup>

• We consider a grid, step  $\Delta$ , • We resort to Taylor series to express  $\Psi(x + \Delta)$  and  $\Psi(x - \Delta)$   $\Psi(x + \Delta) = \Psi(x) + \Delta \frac{d\Psi(x)}{dx} + \frac{\Delta^2 d^2 \Psi(x)}{dx} + \frac{d^2 \Psi(x)}{dx^2} + \frac{d^2 \Psi(x)}{dx^2} + O(\Delta)^5$   $\Psi(x - \Delta) = \Psi(x) - \Delta \frac{d\Psi(x)}{dx^6} + \frac{\Delta^2 d^2 \Psi(x)}{dx^2} + \frac{\Delta^2 d^2 \Psi(x)}{dx^6} + \frac{\Delta^4}{dx} + O(\Delta)^5$ • By summing the above  $\Psi(x + \Delta) = \Psi(x - \Delta) - 2\Psi(x) = \Delta^2 \frac{2 d^2 \Psi(x)}{dx^2} + \frac{\Delta^2 d^2 \Psi(x)}{dx^2} + \Delta^2 \frac{d^2 \Psi(x)}{dx^4} + O(\Delta)^6$ 

 $\Psi(x)$ 

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<sup>e</sup> We consider a grid, step Δ, ⋪ п <sup>e</sup> We resort to Taylor series to express  $\Psi(x + \Delta)$  and  $\Psi(x - \Delta)$   $\Psi(x + \Delta) = \Psi(x) + \Delta \frac{d\Psi(x)}{dx} + \frac{\Delta^2}{\Delta^2} \frac{d\Psi(x)}{dx} + \frac{\Phi^2\Psi(x)}{dx^2} + \frac{\Phi^2\Psi(x)}{dx} + \frac{\Phi^2\Psi($ + O(∆ )<sup>5</sup>  $dx^{2}_{3} d^{6}\Psi(x)$  $\frac{\Delta^2 d^2 \Psi(x)}{\Delta^2 d^2 \Psi(x)} =$  $\Psi(\mathbf{x} - \Delta) = \Psi(\mathbf{x}) - \Delta \frac{\mathrm{d}\Psi(\mathbf{x})}{\mathrm{d}\mathbf{x}^{6^{\dagger}}}$  $\Delta^4 d^4 \Psi(x)$ O(∆ § 24 By summing the above dxexpressions  $\Psi(x - \Delta) - 2\Psi(x) = \Delta$ е  $\frac{dx^2}{dx^2} 2 \frac{d^2 \Psi(x)}{2}$ ₄d4Ψ(x) + O(\( \Delta\))6 dx е We resort to Taylor series to express  $\frac{d_2\Psi(x_2+\Delta)}{d^2\Psi(x_2+\Delta)}$  and  $\frac{d^2\Psi(x_2+\Delta)}{d^2\Psi(x_2+\Delta)} = \frac{d^2\Psi(x_2+\Delta)}{d^2\Psi(x_2+\Delta)}$  $d_2\Psi(\underline{x} - \Delta)$ 2  $O(\Delta^3)$  $\frac{d^2 \Psi(x + \Delta)}{dx} =$  $_{4}\Psi(x$ <u>μ<sup>2</sup>ω2x –Δ</u>  $O(\Delta^3)$ dx dx<sup>4</sup> dx 2 2  $\Delta^{2} d$  $d^2\Psi(x)$  $\Psi(x)$ 

 $\Psi(x)$ 

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<sup>e</sup> We consider a grid, step Δ, ↔ <sup>e</sup> We resort to Taylor series to express  $\Psi(x + \Delta)$  and  $\Psi(x - \Delta)$  $\Psi(\mathbf{x} + \Delta) = \Psi(\mathbf{x}) + \Delta$ + O(∆)<sup>5</sup>  $\Delta_{23}^{2} d^{2} \Psi(x) = \frac{d x^{2}}{\Delta_{3}^{2}} d^{4} \Psi(x)$  $\Psi(\mathbf{x} - \Delta) = \Psi(\mathbf{x}) - \Delta \frac{\mathrm{d}\Psi(\mathbf{x})}{\mathrm{d}\mathbf{x}^{6^{+}}}$  $\frac{\Delta^4}{24} \frac{d^4 \Psi(x)}{du^4}$ O(∆ § expressions  $\Psi(x - \Delta) - 2\Psi(x) = \Delta^{dx^2} 2 \frac{d^{2}\Psi(x)}{d^{2}}$ е ₄d4Ψ(x) + O(\(\Delta\))6 е We resort to Taylor series to express  $\frac{d_2\Psi(x \pm \Delta)}{dx}$  and  $\frac{d_2\Psi(x - \Delta)}{dx}$  $\Delta^{\underline{x}} \Psi(x) + O(\Delta)^3$  $\sqrt{3} d \Psi(x)$  $d^2\Psi(x + \Delta)$ dΦ(x) Δ<sup>4</sup>3Ψ(%)<sup>β</sup> dx d<sup>2</sup>∰(x<sup>2</sup>)  $^{2} \underline{d^{4}\Psi(x)} - O(\Delta^{3})$ d<sup>2</sup>Ψ(x̃−Δ)  $\frac{2}{dx}\frac{4}{dx}\frac{3}{4}$  2 dx<sup>2</sup>  $dx^2$ е  $\lim_{x \to \Delta} \lim_{x \to \Delta} \lim_{x$  $= \Delta^{2} \frac{d^{4}\Psi(x)}{dx^{4}} + O(\Delta)^{5}$ expressions dx dx 2 2 2

 $\Psi(x)$ 

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<sup>e</sup> We consider a grid, step Δ, <sup>e</sup> We resort to Taylor series to express  $\Psi(x + \Delta)$  and  $\Psi(x - \Delta)$   $\Psi(x + \Delta) = \Psi(x) + \Delta \frac{d\Psi(x)}{dx} + \frac{\Delta^2 d\Psi(x)}{2} + \frac{\Phi(X + \Delta)}{2} = \frac{\Phi(X + \Delta)}{6}$ + O(\( \Delta\))5  $\Psi(\mathbf{x} + \Delta) = \Psi(\mathbf{x}) + \Delta \quad \frac{d\mathbf{x}}{d\mathbf{x}} + \frac{d\mathbf{x}}{d\mathbf{x}} +$ <sup>e</sup> By summing the above expressions +  $\Psi(\mathbf{x} - \Delta) - 2\Psi(\mathbf{x}) = \Delta^{dx^2} \frac{2d^2\Psi(\mathbf{x})}{dx^2} + \Delta^{d}\frac{4\Psi(\mathbf{x})}{dx^4} + O(\Delta)^6$ <sup>e</sup> We resort to Taylor series to express  $\frac{d_2\Psi(x \perp \Delta)}{dx^2}$  and  $\frac{d_2\Psi(x \perp \Delta)}{dx^2}$  and  $\frac{d_2\Psi(x \perp \Delta)}{dx^2}$  and  $\frac{d_2\Psi(x \perp \Delta)}{dx^2}$ dx  $\Delta^{2d} {}_{3}\Psi(\mathbf{f}\mathbf{x})^{\beta}$  $\frac{d^2\Psi(x-\Delta)}{dx^2} = \frac{d^2\Psi(x)}{dx^2} - \frac{\Delta^2 d_3\Psi(x)}{2 dx^4 \Delta^4} + 2$  $\frac{2}{d^{4}\Psi(x)} - O(\Delta^{3})$  $B_{\mathcal{A}} \underbrace{\mathsf{Summing}}_{\mathcal{A}} \underbrace{\mathsf{the}}_{\mathcal{A}} \underbrace{\mathsf{above}}_{\mathcal{A}} \underbrace{\mathsf{b}}_{\mathcal{A}} \underbrace{\mathsf{bbove}}_{\mathcal{A}} \underbrace{\mathsf{bb$  $\frac{dx}{dx^{4}} = \Delta^{2} \frac{d^{4}\Psi(x)}{dx^{4}} + O(\Delta)^{5}$ expressions dx e we get  $\Psi(\mathbf{x} + \Delta) + \Psi(\mathbf{x} - \Delta) - 2\Psi(\mathbf{x}) = \frac{\Delta^2}{\frac{1}{2}} \cdot \frac{d^2\Psi(\mathbf{x} + \Delta)}{dx^2} + \frac{d^2\Psi(\mathbf{x} - \Delta)}{\frac{dx}{2}} + 10 \frac{d^2\Psi(\mathbf{x})}{dx^2} + O(\Delta)^6$ 

 $\Psi(x)$ 

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#### The Numerov • Since $\frac{d^2\Psi}{dx} = -Q(x)\Psi(x) + S(x)$ , we get e we get $\sum_{\substack{1 + \frac{\Delta^2}{12} \\ \Sigma}} \sum_{\substack{2 \\ \Sigma}} \sum_{\substack{2 \\ \Sigma}} Q(x + \Delta) \Psi(x + \Delta) = \sum_{\substack{2 \\ \Sigma}}$ $\Psi(x)$ $-1 + \frac{\Delta^2}{12} \frac{Q(x - \Delta) \Psi(x - \Delta)}{\sum_{x \geq 2} \sum_{x \geq 2}}$ $\frac{12}{+2} \sum_{1}^{\Sigma} \frac{5\Delta^2}{12} \frac{\lambda}{Q(x)} \Psi(x)$ а Ь +<sup>2</sup> +12 $(S(x + \Delta) + S(x - \Delta) + 10S(x)) + O(\Delta)$ $\frac{d^{2\psi}}{dx^{2}} + 2 (s - V(x)) \Psi(x) = Q(x) = 2 (s - V(x))$ ))

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S(x) = 0

### The Numerov



$$\sum_{\substack{\lambda = 1}^{\infty}} \sum_{\substack{\lambda = 1 \\ \lambda = 1}}^{\infty} Q(x + \Delta) \frac{\Psi(x + \Delta)}{\Psi(x + \Delta)} = \sum_{\substack{\lambda = 1 \\ \lambda = 1}}^{\infty} \sum_{\substack{\lambda = 1 \\ \lambda = 1}$$

The potential, V (x) is known ;

<sup>e</sup> If we set a value for the total energy of the particule in the box,  $s \rightarrow Q(x) = 2(s - V(x))$  is known

 $+\frac{\Lambda^2}{12}$ 

• the value of the wavefunction at x = a is known :  $\Psi(a) = 0$ ; if we set a value for the wavefunction at  $\overline{\Delta} + \Delta$ , then we get, the value of the wavefunction at  $a + 2\Delta$  $1 + \Delta = Q(a \pm 2\Delta) \frac{\Psi(a + 2\Delta)}{12} = -1 + -Q(a) \frac{\Psi(a)}{12} + 2 = 1 - Q(a \pm \Delta) \frac{\Psi(a + \Delta)}{12}$ 

<sup>e</sup> Then from  $\Psi(a + \Delta)$  and  $\Psi(a + 2\Delta)$ , we can compute  $\Psi(a + 3\Delta)$ , and so on ...

#### Outward integration

### The Numerov



$$\sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma}} Q(x + \Delta) \Psi(x + \Delta) = \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma}} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum_{\substack{\lambda = 1 \\ \Sigma} \sum$$

The potential, V (x) is known ;

<sup>e</sup> If we set a value for the total energy of the particule in the box,  $s \rightarrow Q(x) = 2(s - V(x))$  is known

 $+\frac{\Lambda^2}{12}$ 

• the value of the wavefunction at x = b is known :  $\Psi(b) = 0$  ; if we set a value for the wavefunction at  $\overline{b} - \Delta$ , there we get the value of the wavefunction at  $b - 2\Delta$  $1 + \Delta^2$   $Q(b\Delta^2 2\Delta) \frac{\Psi(b - 2\Delta)}{\Psi(b - 2\Delta)} = -1 + Q(b) \frac{\Psi(b)}{\Psi(b)} + 2 1 - Q(b - \Delta) \frac{\Psi(b - \Delta)}{\Phi(b - \Delta)}$ 

<sup>e</sup> Then from  $\Psi(b - \Delta)$  and  $\Psi(b - 2\Delta)$ , we can compute  $\Psi(b - 3\Delta)$ , and so on ...

#### → Inward integration





Write a code to compute the wavefunctions of the free particule in a box

- 1. Set a guest value for s
- Perform an inward integration from *a* to x<sub>m</sub>, the matching point. The matching point is necessary to get the right value of the energy; in the case of free particle in box problem, a good way to choose the matching point is to take a point where the value of the wavefunctions is different from zero and close to the middle of the box.

3. Perform an outward integration from b to  $x_m$ 

4. Compute the ratios of the first derivative of the wavefunction over the amplitude for both in- and out-ward wavefunctions at the matching point. Change the value of s so that these ratios are identical for both in- and out-ward wavefunctions.

5. Compare the numerical results with the analytical ones.

$$\sum_{\substack{Q(a \pm 2\Delta) \\ Q(a \pm 2\Delta) \\ Q(b \pm 2\Delta) = -1 + \dots \\ Q(a) \\ \Psi(a) + 2 \\ \Psi(a) + 2 \\ \Psi(a) + 2 \\ \Psi(a) \\ \Psi(a) + 2 \\ \Psi(a) \\ \Psi(a) + 2 \\ \Psi(a) \\ \Psi(b) \\$$







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$$-\frac{1}{d^2 \psi x^2} =$$

<sup>e</sup> Set of normalized eigenfunctions

$$\Psi_n(\mathbf{x}) = \sin \frac{2}{\underline{\mu} \pi \underline{x}} , \qquad (1)$$

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where  $n = 1, 2, 3, \cdots$  and L = b - a is the width of the box.

<sup>e</sup> Set of eigenenergies

• Note that 
$$s_1 = 1$$
,  $s_{\frac{\pi}{2}} = \frac{\pi}{2} + s_{\frac{\pi}{2}} = 9$ ,  $\cdots$  if  $L = \frac{\pi}{2}$