## TPMN 2019/2020 : Solving the Schr"odinger

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## TPMN 2019/2020 : Solving the Schr"odinger equation <br> ${ }^{\text {e }}$ H. Bulou, herve.bulou@ipcms.unistra.fr, IPCMS

${ }^{\text {e }}$ Aim of this course : Solving the Sch"odinger equation by using a computer

$$
H \Psi(\mathrm{r})=s \Psi(\mathrm{r})
$$

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$$
H \Psi(\mathrm{r})=\mathrm{s} \Psi(\mathrm{r})
$$

Three kind of terms
${ }^{\mathbf{e}} H$, the Hamiltonian operator $\rightarrow$ it describes the quantum system ; in general it is known
e $\Psi_{(r)}$, the wavefunction of the system ; we want to compute it
$\mathbf{e}_{\mathrm{s} \text {, the total energy of the system ; we want to compute it }}$

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${ }^{\mathbf{e}}$ Problem : the Schr"odinger equation is an eigenvalue problem $\rightarrow$ to get $\Psi(\mathrm{r})$ we need s
we to get s we need $\Psi(\mathrm{r})$ !

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e $\Psi_{(r)}$, the wavefunction of the system; we want to compute it
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Problem : the Schr"odinger equation is an eigenvalue problem $\rightarrow$ to get $\Psi(r)$ we need
we to get s we need $\Psi(\mathrm{r})$ !
From a numerical point of view, there are different ways to solve this problem
${ }^{\mathbf{e}}$ Tomorrow, we will see a quite general method : the Finite Difference Method
${ }^{\mathbf{e}}$ Next week, R. Hertel will present another possible way to proceed : the Finite Element Method

## The Numerov algorithm

${ }^{\mathbf{e}}$ The problem to solve : Free particule in a box (1D)


## The Numerov algorithm

${ }^{\mathbf{e}}$ The problem to solve : Free particule in a box (1D)


$$
-\frac{1 d^{2} \Psi}{2} d x=s \Psi
$$

## The Numerov

 algorithme The problem to solve : Free particule in a box (1D)


$$
-\frac{1 d^{2} \Psi}{2} d x=s \Psi
$$

$$
-\frac{1}{d^{2}} \frac{x^{2}}{2}+V(x) \Psi(x)=
$$

## The Numerov

 algorithme The problem to solve : Free particule in a box (1D)


$$
-\frac{1}{2} \frac{d^{2} \psi}{2} d x=s \psi
$$

$$
-\frac{1}{\frac{1}{\alpha}\left(U_{x}\right.}+V(x) \Psi(x)=
$$

This differential equation belongs to the general kind of $2^{\text {nd }}$ order linear differential equation

$$
\frac{d_{2} \Psi}{d x}+Q(x) \Psi(x)=S(x)
$$

where $Q(x)$ and $S(x)$ are continuous functions on a domain [a, $b]$. The equation is to be solved as a boundary value problem, i.e., $\Psi(a)$ and $\Psi(b)$ are given.

## The Numerov algorithm

$$
\frac{d_{2} \Psi}{d x}+Q(x) \Psi(x)=S(x)
$$

Depending of the functions $Q(x)$ and $S(x)$, the Numerov algorithm can be used to solve
${ }^{\mathbf{e}}$ Eigenvalue problem: $\mathrm{Q}(\mathrm{x}) \boldsymbol{f}=\mathbf{0}$ and $\mathrm{S}(\mathrm{x})=\mathbf{0}$
${ }^{\mathbf{e}}$ The Schr"odinger equation
${ }^{\mathbf{e}}$ Ex. : Hydrogen atom

$Y_{l m}(\theta, \phi)$ are the spherical harmonics and the function $u(r)$ is given by $2^{\text {nd }}$ order differential equation

$$
\frac{d^{2} u}{d r}=-Q(r) u(r) \text { with } Q(r)=2 s+\frac{2 z}{r} \frac{\frac{l(l}{2}}{\frac{r^{2}}{2}}
$$

## The Numerov algorithm

$$
\frac{\mathrm{d} 2}{\mathrm{dx}}+Q(x) \Psi(x)=S(x)
$$

Depending of the functions $Q(x)$ and $S(x)$, the Numerov algorithm can be used to solve

$$
\mathbf{e} \text { Linear system problem: } \mathrm{Q}(\mathrm{x})=\mathbf{0}
$$

The 1D Poisson equation
${ }^{\mathbf{e}}$ Ex. : Hartree potential in spherical symmetry

$$
S(r)=4 \pi r \rho(r)
$$

$$
O(x)=0
$$

$$
\begin{aligned}
& \text { ) }-4 \pi r \rho(r \\
& \text { ) }_{\text {Harte }}(r)=r V_{\text {Harte }}(r \\
& \frac{d 2 U}{\text { dpartree }}=S(r)
\end{aligned}
$$

The Numerov

$$
\frac{d_{2} \Psi}{d x}+Q(x) \Psi(x)=S(x)
$$

algorithm
e The problem to solve : Free particule in a box (1D)


The Numerov

$$
\frac{d 2}{d x}+Q(x) \Psi(x)=S(x)
$$ algorithm

${ }^{\mathbf{e}}$ The problem to solve : Free particule in a box (1D)


$$
\frac{1}{2} \frac{\mathrm{~d}^{2} \Psi}{\mathrm{~d} x^{2}}=\mathrm{s} \Psi
$$

$\frac{d_{2} \Psi}{d x^{2}}+2 s \Psi(x)=0$
$Q(x)=2 s$
$S(x)=0$

The Numerov

$$
\frac{d 2}{d x}+Q(x) \Psi(x)=S(x)
$$

algorithm
${ }^{\mathbf{e}}$ The problem to solve : Free particule in a box (1D)


$$
\frac{1}{2} \frac{d^{2} \psi}{d x^{2}}=s \Psi
$$


$-\frac{1}{s \psi^{2} \psi x^{2}}+V(x) \Psi(x)=$

The Numerov

$$
\frac{d_{2} \Psi}{d x}+Q(x) \Psi(x)=S(x)
$$

algorithm
${ }^{\mathbf{e}}$ The problem to solve : Free particule in a box (1D)


$$
\frac{1}{2} \frac{\mathrm{~d}^{2} \Psi}{\mathrm{dx}}=\mathrm{s} \Psi
$$

$$
\begin{aligned}
& \frac{d_{2} \Psi}{d^{2}}+2 s \Psi(x)=0 \\
& Q(x)=2 s \\
& S(x)=0
\end{aligned}
$$

$$
-\frac{1}{d^{2} \varphi x^{2}}+V(x) \Psi(x)=
$$

$$
\begin{aligned}
& \frac{d^{2} \Psi}{d x}+2(s-V(x)) \Psi(x)=0 \\
& \partial(x)=2(c-V(x
\end{aligned}
$$

$$
Q(x)=2(s-V(x
$$

## The Numerov algorithm

$\mathbf{e}^{\text {We consider a grid, step } \Delta, ~}$


## The Numerov algorithm

${ }^{\mathbf{e}}$ We consider a grid, step $\Delta$,
${ }^{\mathbf{e}}$ We resort to Taylor series to express $\Psi(x+\Delta)$ and $\Psi(x-\Delta)$

$$
\begin{aligned}
& \Psi(x+\Delta)=\Psi(x)+\Delta \quad \frac{d \psi(x)}{d x} \quad \Lambda_{2}^{d} \frac{d^{2} \psi(x)}{} f^{6} \frac{U x)}{d x^{2}} \frac{d^{4} \Psi(x)}{6}-1-O(\Delta)^{5}
\end{aligned}
$$

$$
\mathrm{d} x^{2} \quad \mathrm{~d} x^{6}
$$

## The Numerov algorithm

${ }^{\mathbf{e}}$ We consider a grid, step $\Delta$,
${ }^{\mathbf{e}}$ We resort to Taylor series to express $\Psi(x+\Delta)$ and $\Psi(x-\Delta)$
${ }^{\mathbf{e}}$ By summing the above


## The Numerov algorithm

${ }^{\mathbf{e}}$ We consider a grid, step $\Delta$,
${ }^{\mathbf{e}}$ We resort to Taylor series to express $\Psi(x+\Delta)$ and $\Psi(x-\Delta)$

${ }^{\mathbf{e}}$ By summing the above expf(XStoA $s^{+} \Psi(x-\Delta)-2 \Psi(x)=\Delta^{\frac{d x^{2}}{2}} \frac{2 d^{2} \psi(x)}{d x^{2}}+\frac{A^{4} L^{4} \Psi(x)}{d x^{4}}+O(\Delta)^{6}$


$$
\begin{aligned}
& d^{2} \omega(x+\Delta) \\
& \frac{d}{2} \frac{d}{d x}\left(\frac{d}{d x}=1\right. \\
& \frac{2}{2}=
\end{aligned}
$$



## The Numerov algorithm

${ }^{\mathbf{e}}$ We consider a grid, step $\Delta$,
e We resort to Taylor series to express $\Psi(x+\Delta)$ and $\Psi(x-\Delta)$

${ }^{\mathbf{e}}$ By summing the above




## The Numerov algorithm

e We consider a grid, step $\Delta$,
${ }^{\mathbf{e}}$ We resort to Taylor series to express $\Psi(x+\Delta)$ and $\Psi(x-\Delta)$


$$
\begin{array}{ll}
\Psi(x+\Delta)=\Psi(x)+\Delta & \frac{d \Psi(x)}{d x} \\
\Psi(x-\Delta)=\Psi(x)-\Delta \frac{\Delta^{2}}{2} \frac{d^{2} \Psi(x)}{d x^{6}} & \Delta^{2} d^{4} d^{4} \Psi(x) \\
d^{4} x^{4} & -\frac{\Delta^{3}}{6} \frac{d^{6} \Psi(x)}{d x^{6}}+\frac{\Delta^{4}}{24} \frac{d^{4} \Psi(x)}{d x^{4}}-O(\Delta)^{5}
\end{array}
$$

${ }^{\mathbf{e}}$ By summing the above expHesstôs $+\Psi(x-\Delta)-2 \Psi(x)=\Delta^{d x^{2}} \frac{d^{2} \Psi(x)}{d x^{2}}+\Delta_{12}^{d^{4}} \frac{d(x)}{d x^{4}}+O(\Delta)^{6}$
${ }^{\mathbf{e}}$ We resort to Taylor series to express ${ }^{d_{2} \psi \partial_{x}(\Delta)}$ and $d_{2} \psi(x-\Delta)$

e we get

## The Numerov

 algorithm
e Since $\frac{d}{d} \frac{2}{2} \underline{2}=-Q(x) \Psi(x)+S(x)$, we get
e we get ${ }^{2}$


$$
\begin{array}{cc}
\Sigma & \Sigma \\
1+\frac{\Delta^{2}}{12} & Q(x+\Delta) \Psi(x+\Delta)= \\
\Sigma & \Sigma \\
-1+\frac{\Delta^{2}}{12} Q(x-\Delta) \Psi(x-\Delta) \\
\Sigma & +21-\frac{5 \Delta^{2}}{12} Q(x) \Psi(x)
\end{array}
$$

$+\frac{t^{2}}{12}$

$$
(S(x+\Delta)+S(x-\Delta)+10 S(x))+\mathrm{O}_{( }^{6}(\Delta)
$$

$$
\begin{aligned}
& \frac{d^{2} \psi}{d x^{2}}+2(s-V(x)) \Psi(x)= \\
& Q(x \quad \theta=2(s-V(x \\
& )) \\
& S(x)=0
\end{aligned}
$$

## The Numerov



$$
\begin{aligned}
& \frac{d^{2} \Psi}{d x^{2}}+2(s-V(x)) \Psi(x)= \\
& d^{2}(x \emptyset=2(s-V(x \\
& S(x)=0
\end{aligned}
$$

$$
\begin{aligned}
& \Sigma \quad \Sigma \\
& \begin{array}{c}
1+\frac{\Delta^{2}}{12}
\end{array} Q(x+\Delta) \Psi(x+\Delta)= \\
& -1+\frac{\Delta^{2}}{12 \sum \Sigma(x-\Delta) \Psi(x-\Delta)} \\
& +21-\frac{5 \Delta^{2}}{12^{Q}(x)}(x) \\
& +\frac{\Delta^{2}}{12}(S(x+\Delta)+S(x-\Delta)+10 S(x))+O\left(\delta^{\delta}\right.
\end{aligned}
$$

The potential, $V(x)$ is known ;
e If we set a value for the total energy of the particule in the box, $s \rightarrow Q(x)=2(s-V(x))$ is known
e the value of the wavefunction at $x=a$ is known: $\Psi(a)=0$; if we set a value for the wavefunction

$1+\frac{\Delta}{12} \quad Q\left(a \AA^{2} 2 \Delta\right) \Psi(a+2 \Delta)=-1+\longrightarrow(a) \Psi(a)+21-\quad 12(a+\Delta) \Psi(a+\Delta)$
Then from $\Psi(a+\Delta)$ and $\Psi(a+2 \Delta)$, we can compute $\Psi(a+3 \Delta)$, and so on $\ldots$

## The Numerov



$$
\begin{aligned}
& \frac{d^{2} \Psi}{d x^{2}}+2(s-V(x)) \Psi(x)= \\
& d(x \emptyset=2(s-V(x \\
& S(x)=0
\end{aligned}
$$

$$
\begin{aligned}
& \Sigma \quad \Sigma \\
& \begin{array}{c}
1+\frac{\Delta^{2}}{12}
\end{array} Q(x+\Delta) \Psi(x+\Delta)= \\
& -1+\frac{\Delta^{2}}{12 \sum \Sigma(x-\Delta) \Psi(x-\Delta)} \\
& +21-\frac{5 \Delta^{2}}{12^{Q}(x)} \Psi(x) \\
& { }^{+} \frac{\Delta^{2}}{12}(S(x+\Delta)+S(x-\Delta)+10 S(x))+O(\wp
\end{aligned}
$$

The potential, $V(x)$ is known ;
e If we set a value for the total energy of the particule in the box, $s \rightarrow Q(x)=2(s-V(x))$ is known
e the value of the wavefunction at $x=b$ is known: $\Psi(b)=0$; if we set a value for the wavefunction at $\Sigma-\Delta$, then $\sum^{2}$ we get the value of the wavefunctio $\sum_{1}$ at $b-2 \Sigma^{\Sigma}$
e Then from $\Psi(b-\Delta)$ and $\Psi(b-2 \Delta)$, we can compute $\Psi(b-3 \Delta)$, and so on $\ldots$
The Numerov


$$
\begin{aligned}
& \frac{d^{2} \mu}{d x^{2}}+2(s-V(x)) \Psi(x)= \\
& Q(x \emptyset=2(s-V(x \\
& S(x)=0
\end{aligned}
$$



Write a code to compute the wavefunctions of the free particule in a box

1. Set a guest value for $s$
2. Perform an inward integration from $a$ to $x_{m}$, the matching point. The matching point is necessary to get the right value of the energy ; in the case of free particle in box problem, a good way to choose the matching point is to take a point where the value of the wavefunctions is different from zero and close to the middle of the box.
3. Perform an outward integration from $b$ to $x_{m}$
4. Compute the ratios of the first derivative of the wavefunction over the amplitude for both in- and out-ward wavefunctions at the matching point. Change the value of $s$ so that these ratios are identical for both in- and out-ward wavefunctions.
5. Compare the numerical results with the analytical ones.


$$
\begin{aligned}
& \sum_{Q(a)}^{\Sigma} \underset{(a)+2}{\Sigma} \underset{=}{y_{\Delta}^{2}} \\
& \Sigma \\
& Q(a+\Delta) \Psi(a \\
& Q(b-\Delta) \Psi(b
\end{aligned}
$$

## The Numerov algorithm





## The Numerov algorithm








## The Numerov


${ }^{\text {e }}$ Set of normalized
eigenfunctions

$$
\begin{equation*}
\Psi_{n}(x)=\sin _{\frac{2}{n \pi x}}^{\Sigma} \tag{1}
\end{equation*}
$$

where $n=1,2,3, \cdots$ and $L=b-a$ is the width of the box.
${ }^{\mathbf{e}}$ Set of eigenenergies

$$
\begin{equation*}
2 s_{n} \frac{\Pi^{2} n}{2 L^{2}} \tag{2}
\end{equation*}
$$

e Note that $\mathrm{s}_{1}=1, \mathrm{~s} \underset{3}{\bar{z}} 4, \mathrm{~s}=9, \cdots$ if $L=\frac{\sqrt{\pi}}{2}$

